

HYDRAULIC FRACTURE

Hydraulic Fracture with Leak off

INTRODUCTION

- Hydraulic fracturing is a process that is used to create fractures in rocks.
- It was first used in the US in 1947, and went commercial in 1949.
- Its success in increasing production from oil wells made it to be adapted worldwide.
- The most important industrial use is in stimulating oil and gas wells.

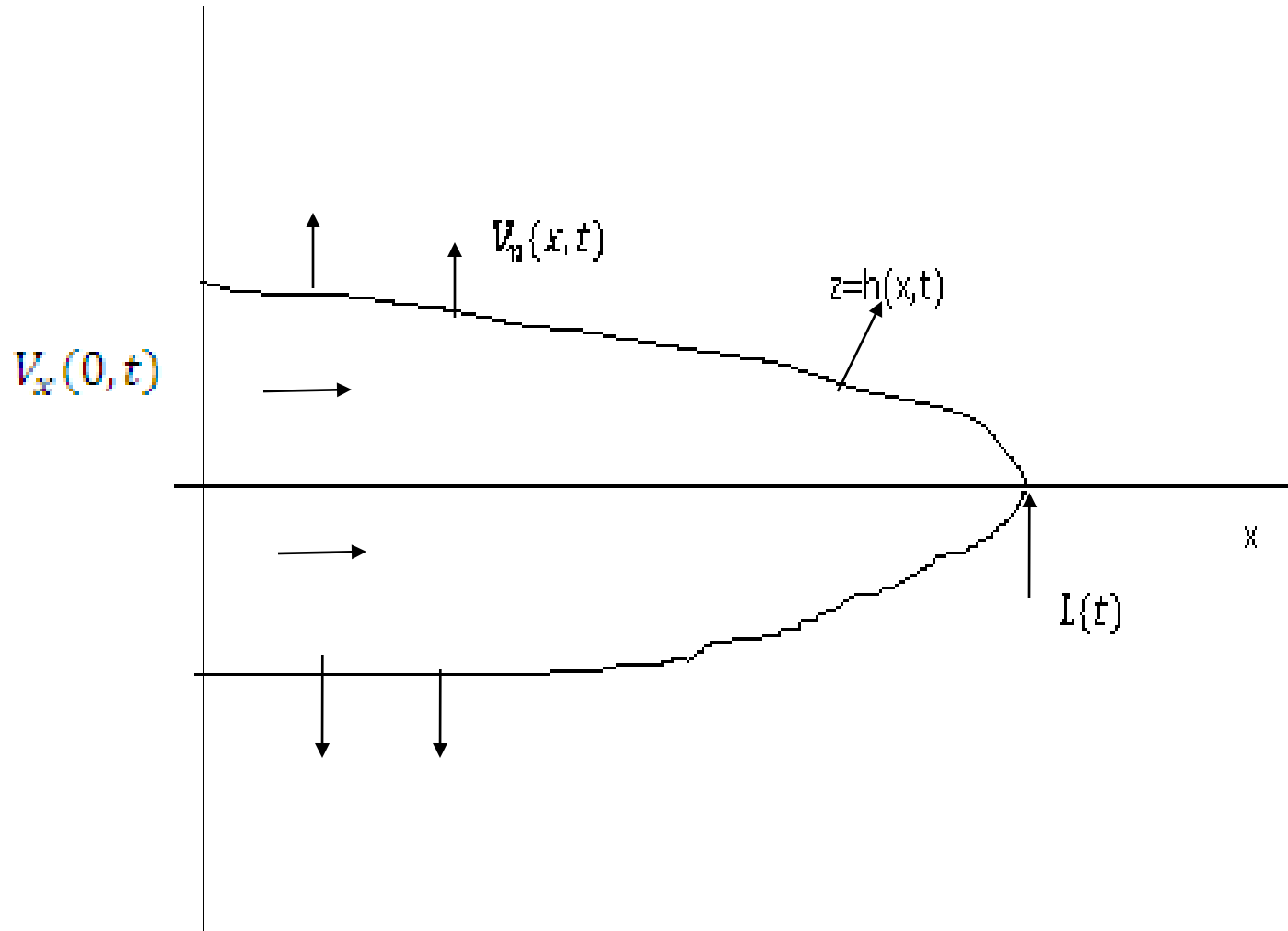
- Investigate effect of leak-off on growth of fracture
- Investigate the extraction of fluid from fracture.

Hydraulic Fracture with Leak off

- The fluid is pumped into the fracture by a fluid injection at a velocity v_x
- The cavity walls are permeable so some amount of fluid escapes into or its sucked out of the permeable rock at a leak off velocity

$$v_n(x, t)$$

- $h(x, t) = \frac{1}{2} \times \text{width of fracture.}$



- The PDE that describes the situation mathematically is given by:

$$\frac{\partial h(x, t)}{\partial t} = \frac{\Lambda}{3\mu} \frac{\partial}{\partial x} \left(h^3 \frac{\partial h(x, t)}{\partial x} \right) - v_n(x, t)$$

- Firstly, we derive the above equation
- Then the above PDE was solved analytically using scaling transformation, similarity solutions.

CHARACTERISTIC EQUATIONS

$$\bar{V}_x = \frac{V_x}{U}, \quad \bar{p} = \frac{p}{P}, \quad P = \frac{\mu LU}{H^2}$$

$$\bar{V}_z = \frac{V_z}{W}, \quad Re = \frac{UL}{\nu}, \quad \nu = \frac{\mu}{\rho}$$

BY SUBSTITUTING THE ABOVE SCALAR EQUATIONS INTO EULERS EQUATION WE GET

$$Re \left(\frac{H}{L} \right)^2 \left[\frac{\partial \bar{V}_x}{\partial t} + \bar{V}_x \frac{\partial \bar{V}_x}{\partial \bar{x}} + \bar{V}_z \frac{\partial \bar{V}_x}{\partial \bar{z}} \right] = - \frac{\partial \bar{p}}{\partial \bar{x}} + \left(\frac{H}{L} \right)^2 \frac{\partial \bar{V}_x}{\partial \bar{x}^2} + \frac{\partial^2 \bar{V}_x}{\partial \bar{z}^2} \dots \dots \dots (1)$$

$$Re \left(\frac{H}{L} \right)^2 \left[\frac{\partial \bar{V}_z}{\partial t} + \bar{V}_x \frac{\partial \bar{V}_z}{\partial \bar{x}} + \bar{V}_z \frac{\partial \bar{V}_z}{\partial \bar{z}} \right] = - \frac{\partial \bar{p}}{\partial \bar{z}} + \left(\frac{H}{L} \right)^4 \frac{\partial \bar{V}_z}{\partial \bar{x}^2} + \left(\frac{H}{L} \right)^2 \frac{\partial^2 \bar{V}_z}{\partial \bar{z}^2} \dots \dots \dots (2)$$

CONTINUITY EQUATION

$$\frac{\partial \bar{V}_x}{\partial \bar{x}} + \frac{\partial \bar{V}_z}{\partial \bar{z}} = 0 \dots \dots \dots (3)$$

LUBRICATION APPROXIMATION

$$\frac{H}{L} \ll 1, \quad Re \left(\frac{H}{L} \right)^2 \ll 1$$

$$\frac{\partial^2 \bar{V}_x}{\partial \bar{z}^2} = \frac{\partial \bar{p}}{\partial \bar{x}}$$

$$\frac{\partial \bar{p}}{\partial \bar{z}} = 0$$

$$\frac{\partial \bar{V}_x}{\partial \bar{x}} + \frac{\partial \bar{V}_z}{\partial \bar{z}} = 0$$

Boundary Conditions

$$z = h(x, t)$$

$$V_x(x, h, t) = 0$$

$$z = h(x, t)$$

$$V_z(x, h, t) = \frac{\partial h}{\partial t} + V_n(x, t)$$

$$z = -h(x, t)$$

$$V_x(x, -h, t) = 0$$

$$z = -h(x, t)$$

$$V_z(x, -h, t) = -\frac{\partial h}{\partial t} - V_n(x, t)$$

Using the continuity equation (3), we first showed that

$$\frac{\partial h}{\partial t} + \frac{1}{2} \frac{\partial}{\partial x} \int_{-h(x,t)}^{h(x,t)} V_x(x, z, t) dz = -V_n(x, t) \dots \dots \dots (4)$$

We solved for $V_x(x, z, t)$

$$V_x(x, z, t) = \frac{1}{2\mu} \frac{\partial p(x, t)}{\partial x} (z^2 - h^2)$$

Substituting into (4),

$$\frac{\partial h}{\partial t} = \frac{1}{3\mu} \frac{\partial}{\partial x} \left(h^3 \frac{\partial p}{\partial x} \right) - V_n(x, t)$$

Using simple model, we took $p(x, t) = \Lambda h(x, t)$

Thus we have our equation

$$\frac{\partial h}{\partial t} = \frac{\Lambda}{3\mu} \frac{\partial}{\partial x} \left[h^3 \frac{\partial h}{\partial x} \right] - V_n(x, t)$$

PDE

$$\frac{\partial h}{\partial t} = \frac{\lambda}{3\mu} \frac{\partial}{\partial x} \left[h^3 \frac{\partial h}{\partial x} \right] - V_n(x, t)$$

WE USE THE FOLLOWING FUNCTIONS:

$$h(x, t) = t^{\frac{1}{3}(2\alpha-1)} F(\xi)$$

$$V_n(x, t) = t^{\frac{1}{3}(\alpha-2)} G(\xi)$$

WHERE: $\xi = \frac{x}{t^\alpha}$

$$\frac{\lambda}{3\mu} \frac{d}{d\xi} \left[F^3(\xi) \frac{dF(\xi)}{d\xi} \right] + \alpha \frac{d}{d\xi} (\xi F) - \frac{5}{3} \left(\alpha - \frac{1}{5} \right) F(\xi) - G(\xi)$$

BY USING

$$H(\eta) = AF(\xi), \quad F(\xi) = \frac{1}{A}H(\eta)$$

$$W(\eta) = BG(\xi), \quad G(\xi) = \frac{1}{B}W(\eta)$$

$$\text{WHERE: } \eta = \frac{\xi}{L_0} = \frac{x}{L_0 t^\alpha} = \frac{x}{L(t)}$$

$$h(x, t) = t^{\frac{1}{3}(2\alpha-1)} \left(\frac{3\mu L_0^2}{\lambda} \right)^3 H(\eta)$$

WE END UP WITH

$$\frac{d}{d\eta} \left[H^3 \frac{dH}{d\eta} \right] + \alpha \frac{d}{d\eta} (\eta H) - \frac{5}{3} \left(\alpha - \frac{1}{5} \right) - W = 0$$

We now have the following ODE

$$\frac{d}{dy} \left[H^3 \frac{d}{dy} H \right] + \alpha \frac{d}{dy} [\gamma H] - \left(\frac{5}{3} \left(\alpha - \frac{1}{5} \right) \right) H - W = 0$$

with boundary condition $H(1) = 0$

Assumption: $W = \beta H$

$$\frac{d}{dy} \left[H^3 \frac{d}{dy} H \right] + \alpha \frac{d}{dy} [\gamma H] - \left(\frac{5}{3} \left(\alpha - \frac{1}{5} \right) \right) H - \beta H = 0$$

Case 1: $\beta = \frac{5}{3} \left(\frac{1}{5} - \alpha \right)$

Substituting in we obtain:

$$\frac{d}{d\gamma} \left[H^3 \frac{d}{d\gamma} H \right] + \alpha \frac{d}{d\gamma} [\gamma H] = 0, H(1) = 0$$

solving this we obtained the following results:

$$H(\gamma) = \left[\frac{3}{2} \alpha (1 - \gamma^2) \right]^{\frac{1}{3}}$$

$$h(x, t) = t^{\frac{1}{3}(2\alpha-1)} \left(\frac{\Lambda L_o^2}{3\mu} \right)^{\frac{1}{3}} H \left(\frac{x}{L_o t^\alpha} \right)$$

Interpretation:

volume in cavity $V(t) = 2 \int_0^{L(t)} h(t, x) dx$

$$\frac{dV}{dt} = \int_{-h}^h V_x(0, t) dz - 2 \int_0^{L(t)} V_n(t, x) dx$$

Constant volume : $\alpha = \frac{1}{5}$

dV/dt constant: $\alpha = \frac{4}{5}$

Constant pressure at x = 0: $\alpha = \frac{1}{2}$

Case 2:

We left the ODE as is and made the ansatz

$$H(\gamma) = A[1 - \gamma]^n$$

Substituting in, we obtained

$$n = \frac{1}{3} \quad A = (3\alpha)^{1/3}$$

Solving for h we obtained

$$h(x, t) = t^{\frac{1}{3}(2\alpha-1)} \left(\frac{\Lambda L_o^2}{3\mu} \right)^{\frac{1}{3}} H \left(\frac{x}{L_o t^\alpha} \right)$$

The relationship between α and β can be expressed as

$$\alpha = 1 - 3\beta$$

For $\beta = 0$, $L(t) = L_0 t$ where $L(t)$ is the length of the cavity and L_0 is some constant.

Conditions on α and β for inflow and outflow at the entrance of the cavity:

$$Q_1 = \int_{-h(0,t)}^{h(0,t)} V_x(0, z, t) dz$$

$$Q_2 = 2 \int_0^{L(t)} V_n(t, x) dx$$

$$V(t) = 2 \int_0^{L(t)} h(t, x) dx$$

$$\frac{dV}{dt} + Q_2 = Q_1$$

Results:

$$\beta + \frac{5}{3} \left(\alpha - \frac{1}{5} \right) = 0$$

no injection of fluid

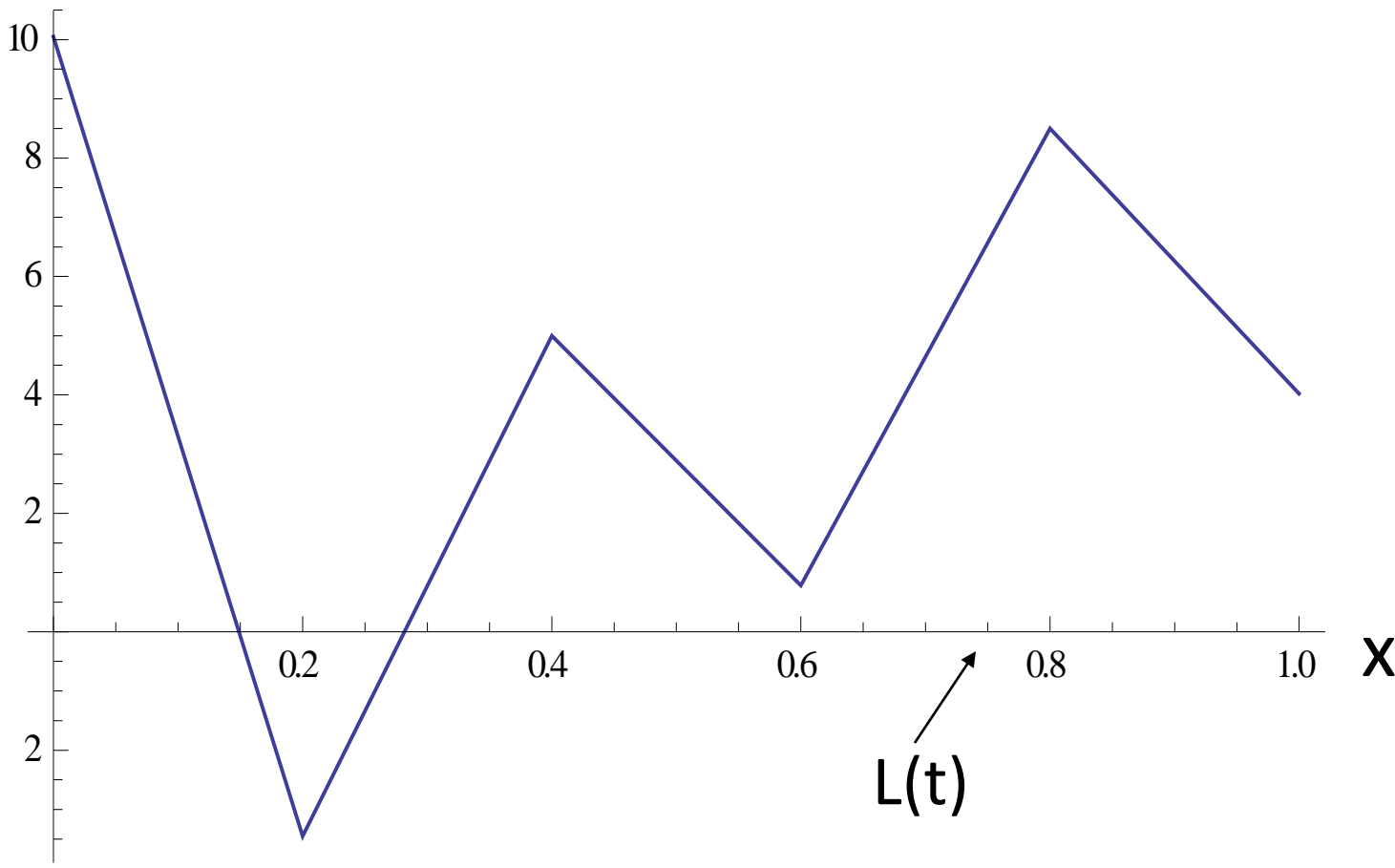
$$\beta + \frac{5}{3} \left(\alpha - \frac{1}{5} \right) > 0$$

fluid is injected

$$\beta + \frac{5}{3} \left(\alpha - \frac{1}{5} \right) < 0$$

fluid is extracted

A noisy solution due to the discontinuity



Finite Difference Method

$n = \text{number of nodal points along } \gamma \text{ axis}$

$\Delta = \text{step length}$

$H_i \approx H(\gamma_i)$

$$\frac{dH}{d\gamma}(\gamma_i) = \frac{-H_{i-1} + H_{i+1}}{\Delta} + O(\Delta)$$

$$\frac{d^2H}{d\gamma^2}(\gamma_i) = \frac{H_{i-1} - 2H_i + H_{i+1}}{\Delta^2} + O(\Delta^2)$$

Method of Lines

$$\frac{\partial h}{\partial t} = \frac{\Lambda}{3\mu} \frac{\partial}{\partial x} \left[h^3 \frac{\partial h}{\partial x} \right]$$

n = number of nodal points along γ axis

Δ = step length

$h_i(t) \approx h(x_i, t)$

$$\frac{\partial h}{\partial x}(t, x_i) = \frac{-h_{i-1}(t) + h_{i+1}(t)}{\Delta} + O(\Delta)$$

$$\frac{\partial^2 h}{\partial x^2}(t, x_i) = \frac{h_{i-1}(t) - 2h_i(t) + h_{i+1}(t)}{\Delta^2} + O(\Delta^2)$$

- Plot of numerical and analytical solutions

